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## THEORETICAL STUDIES FOR THIRD-ORDER HYPERPOLARIZABILITIES OF ALTERNANT AND CONDENSED-RING CONJUGATED SYSTEMS I

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**Abstract** Static third-order hyperpolarizabilities( $\gamma$ ) for neutral, ion radical and ion states of centrosymmetric systems are calculated by the use of the *ab initio* method. The dependences of basis sets and electron correlation effects on the  $\gamma$  values for allyl systems are examined. It is found that the extended basis set with diffuse *p* and *d* functions and Møller-Plesset second-order perturbation (MP2) method are necessary for reasonable descriptions of at least relative tendencies for these charged states of allyl. We also show that triafulvalene (alternant conjugated system) exhibits positive  $\gamma$  value, while the anion radical state of pentalene (condensed-ring conjugated system) does negative one by the *ab initio* MP2 calculations using 6-31G+*pd*( $\zeta_{p,d}=0.05$ ).

### INTRODUCTION

Recently, considerable attention has been given to third-order nonlinear optical effects on organic substances. Many theoretical calculations of  $\gamma$ , which is the origin of macroscopic third-order nonlinear optical effects, are carried out.<sup>1-14</sup> From the time-dependent perturbational analysis of nonlinear optical processes, it is well known that the  $\pi$ -conjugated noncentrosymmetric systems tend to possess large positive  $\gamma$  values. Previously, we presented the three-type approximation<sup>1,3,5</sup> in which virtual excitation processes describing the  $\gamma$  can be divided into three types: (I) ( $\gamma > 0$ ), (II)' ( $\gamma < 0$ ) and (III)' ( $\gamma > 0$ ). The three-type approximation shows that the signs and magnitudes of  $\gamma$  for centrosymmetric systems are determined from the detailed balance between types (II)' and (III)'. We reformulated the *N* state three-type approximate formula by the use of the ground and excited state polarizabilities and the ratio of excitation energies as follows.<sup>1</sup>

$$\gamma^{(I)+(II)'+(III)} = \frac{1}{2} \sum_{n=1}^N \left\{ \frac{\alpha_0(n)}{E_{n0}} \left[ \frac{(\Delta\mu_{n0})^2}{E_{n0}} + \frac{1}{2} \sum_{m(\neq n)=1}^N \left( \alpha_n(m)r(m,n) - [1-r(m,n)] \frac{\alpha_0(n)}{N-1} \right) \right] \right\}. \quad (1)$$

Here,  $\alpha_0(n)$  is the contribution of the  $n$ th excited states to the ground state polarizability  $\alpha_0$  and  $\alpha_n(m)$  is the contribution of the  $m$ th excited states to the  $n$ th excited state polarizability  $\alpha_n$ . The factor  $r(m,n)$  is the ratio of excitation energy  $E_{mn}$  to  $E_{n0}$ . The first and second terms of Eq. (1) correspond to the type (I) and type (II+III)', respectively. The centrosymmetric system (type (I)=0), whose  $N$  value is large and whose  $r(m,n)$  is not close to 0 (nondegenerate), tends to exhibit a positive  $\gamma$  value, whereas that with numerous near-degenerate excited states has a possibility of exhibiting not only positive but also negative  $\gamma$  values. Furthermore, in the limiting case with the smallest  $N$  value ( $N=3$ ), Eq. (1) can be reduced to the three-state three-type approximate formula;<sup>1</sup>

$$\gamma^{(I)+(II)+(III)} = \frac{1}{2} \frac{\alpha_0}{E_{n0}} \left( \frac{(\Delta\mu_{n0})^2}{E_{n0}} + \frac{1}{2} [\alpha_n r - \alpha_0 (1-r)] \right), \quad (2)$$

where the first term represents the type (I) and the second term does to the type (II+III)'. The formula shows that the centrosymmetric three state systems have negative  $\gamma$  values in the case of  $[\alpha_n > 0 \text{ and } \alpha_n / \alpha_0 < E_{n0} / E_{mn}]$  or  $\alpha_n < 0$ . The condition seems to be feasible to be achieved for the centrosymmetric systems with large ground state polarizability  $\alpha_0 (> \alpha_n)$ . Therefore, we investigate three charged states (neutral, ion radical and ion states) of centrosymmetric alternant and condensed-ring conjugated model systems with large polarizability  $\alpha_0$ .

First, we have to check the basis set and electron correlation dependences of the  $\gamma$  values in order to choose appropriate size basis sets and correlation methods before calculating the  $\gamma$  values of alternant and condensed-ring systems. In this end, the  $\gamma$  values of neutral, cation and anion states of allyl are examined with respect to the various basis sets and electron correlation methods. The  $\gamma$  values are calculated by the finite-field (FF) method using the coupled Hartree-Fock (CHF) + electron correlation (MPn, coupled-cluster(CC)) methods.<sup>15</sup> In the first step, the restricted HF (RHF) and the unrestricted HF (UHF) methods are used for closed and open shell systems, respectively. In this study, the *ab initio* calculations are performed by using GAUSSIAN 92 program package.<sup>16</sup>

Second, as models of alternant and condensed-ring conjugated systems, main components of  $\gamma$  values for neutral, ion radical and ion states of triafulvalene and pentalene are calculated by the methods selected above. We discuss the results in relation to various charged states, molecular structures and correlation effects.

$\gamma$  VALUES OF NEUTRAL, CATION AND ANION STATES OF ALLYL

We investigate main components of  $\chi(\gamma_{xxx})$  for neutral (a), cation (b) and anion (c) states of allyl shown in Figure 1. In the case, according to the three-type approximation, the systems have a possibility to exhibit negative  $\gamma_{xxx}$  values except positive ones. The geometry of each charged state of allyl is fully optimized by the *ab initio* HF calculations using 6-31G\*\* basis set. The basis sets used for calculations of  $\gamma_{xxx}$  are STO-3G, 6-31G, 6-31G\*\*, 6-31+G ( $\zeta_{sp}=0.0438$ ), 6-31G+d ( $\zeta_d=0.05$ ), 6-31G+pd ( $\zeta_p=0.05$ ,  $\zeta_d=0.05$ ) and 6-31G+pdd ( $\zeta_p=0.05$ ,  $\zeta_d=0.2$ ,  $\zeta_d=0.05$ ). We also examine the effects of electron correlations on  $\gamma_{xxx}$  values. The applied correlation methods are MP2, MP3, MP4DQ(including double(D) and quadruple(Q) excitations), MP4SDQ(including single(S), double(D) and quadruple(Q) excitations), MP4SDTQ(including single(S), double(D), triple(T) and quadruple(Q) excitations), CCSD(including single(S) and double(D) excitations) and CCSD(T) (including single(S), double(D) and triple(T) excitations). We use 0.01 a.u. for neutral and cation states, while use 0.005 a.u. for anion states as the minimum electric field ( $F^*$ ) in the FF method. The  $\gamma_{xxx}$  value is obtained by the numerical differentiation of total energy  $E$  with respect to the field  $F^*$  as follows.

$$\gamma_{xxx} = -\{56E(0) - 39[E(F^*) + E(-F^*)] + 12[E(2F^*) + E(-2F^*)] - [E(3F^*) + E(-3F^*)]\} / 6(F^*)^4 \quad (3)$$

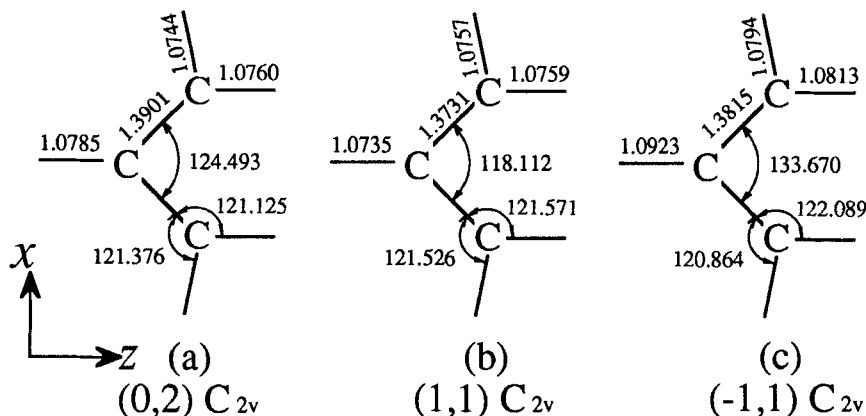


FIGURE 1 Structures of neutral (a), cation (b) and anion (c) allyl systems optimized by *ab initio* HF method using 6-31G\*\*. Bond lengths [Å] and bond angles [°] are shown. Numbers in parentheses depict charge and spin multiplicity, respectively.

Figure 2 shows variations of  $\gamma_{xxx}$  of neutral (a), cation (b) and anion (c) states of allyl with basis sets and electron correlations. First, as for the neutral allyl (a), the basis sets excluding diffuse  $p$  and  $d$  functions (STO-3G, 6-31G and 6-31G\*\*) are found to give small  $|\gamma_{xxx}|$  and exhibit small electron correlation effects on  $\gamma$  values. In contrast, extended basis sets including diffuse  $p$  and/or  $d$  functions (6-31+G, 6-31G+ $d$ , 6-31G+ $pd$  and 6-31G+ $pdd$ ) provide large  $|\gamma_{xxx}|$  and exhibit similar tendencies of correlation effects with each other. Therefore, it is concluded that variations of  $|\gamma|$  with electron correlation effects cannot be reproduced correctly using the basis sets without diffuse  $p$  and  $d$  functions. It is also found that the  $|\gamma_{xxx}|$  values increase with the size of basis set till the 6-31G+ $pd$  ( $\zeta_p=0.05$ ,  $\zeta_d=0.05$ ), while the 6-31G+ $pdd$  in which the tight  $d$  ( $\zeta_d=0.2$ ) function is added to the 6-31G+ $pd$ , tends to fall the  $|\gamma_{xxx}|$  values slightly. The MP2 correlation effects on  $\gamma$  values are found to be largest and be positive in sign. The change of  $\gamma_{xxx}$  from MP2 to MP4 is found to be small. The effects of S, D and Q excitations decrease the  $|\gamma_{xxx}|$ , while those of T increase the  $|\gamma_{xxx}|$  slightly. The CCSD method, in which the S and D excitation effects are included until the infinite order, is found to give considerably smaller  $|\gamma_{xxx}|$  than those of MP4. This suggests that higher-order (beyond the fourth-order) correlation corrections are needed to evaluate the  $|\gamma_{xxx}|$  quantitatively. The  $|\gamma_{xxx}|$  values by CCSD(T) method are shown to be slightly increased against the CCSD method due to the T excitation effects.

Second, the basis set and electron correlation dependences of  $\gamma_{xxx}$  for allyl cation (b) are examined. As shown in Figure 2, the  $\gamma_{xxx}$  values by all basis sets and correlation methods are negative in sign in contrast to those of neutral allyl (a). The behaviors of  $\gamma_{xxx}$  with the correlation effects are all similar regardless of their used basis sets. The differences in the magnitudes of  $\gamma_{xxx}$  and the tendencies of variations of  $\gamma_{xxx}$  with the correlation effects are shown to be smaller than those of neutral allyl (a). This seems to be caused by the unnecessary of diffuse  $p$  and  $d$  functions in describing cation states. On the other hand, the changes of  $\gamma_{xxx}$  of allyl cation (b) with electron correlations are much larger than those of neutral allyl (a). It is revealed that the correlation effects contribute to the increase of  $|\gamma_{xxx}|$  from MP2 to MP4DQ, while S and T excitation effects in MP4 contribute to the decrease of  $|\gamma_{xxx}|$ . It is also found that the  $\gamma_{xxx}$  values of CCSD are close to the MP4SDQ values and the CCSD(T) tends to decrease the  $|\gamma_{xxx}|$  due to the T excitation effects. For allyl cation (b), although each order correlation effect is considerably large, the CCSD(T) values are shown to be close to the MP2 values eventually.

Third, the  $\gamma_{xxx}$  values of allyl anion (c) are investigated. As can be seen from the variations of  $\gamma_{xxx}$  shown in Figure 2, the basis set dependence is largest among three

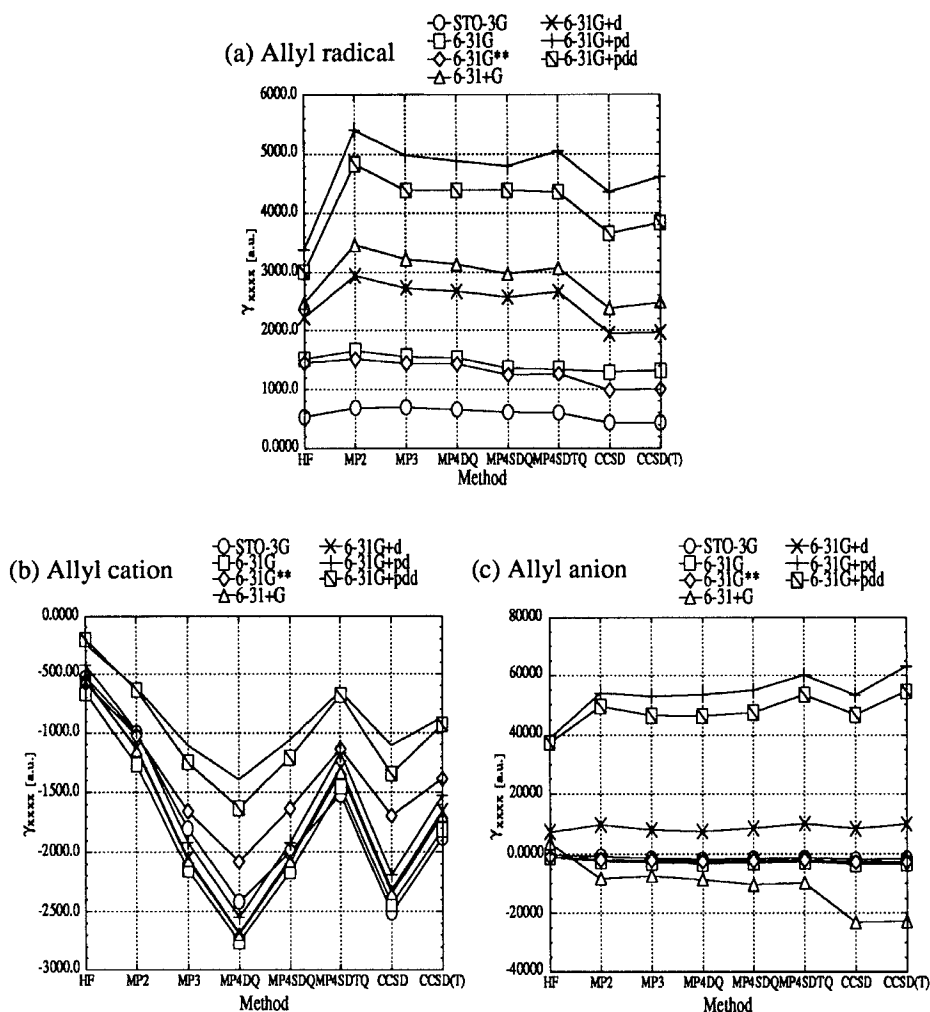


FIGURE 2 Basis set and electron correlation dependences of  $\gamma_{xxxx}$  for neutral (a), cation (b) and anion (c) states of allyl systems.

charged states (a)-(c). This seems to be a result predicted from the necessity of diffuse  $p$  and/or  $d$  functions in describing anion states. It is found that the  $\gamma_{xxxx}$  values using the basis sets without both diffuse  $p$  and  $d$  functions come to be small negative values and the variations of  $\gamma_{xxxx}$  with correlation effects are also small. Although the 6-31+G basis set is usually available for describing anion states, that is found to be unavailable for

reproducing at least the tendencies of  $\gamma$  values of anion states: the  $\gamma_{xxx}$  value using 6-31+G at the HF level is shown to give the positive small value, whereas higher-order correlation methods fall off the  $\gamma_{xxx}$  to negative values. On the other hand, all correlation methods using the 6-31G+d are shown to provide the positive  $\gamma_{xxx}$  values, but the magnitudes of  $\gamma_{xxx}$  and the correlation effects are relatively small. These results indicate that the extended basis sets only including either diffuse *p* or *d* function cannot provide reasonable tendencies of  $\gamma$  for anion states. The methods using 6-31G+*pd* basis set are found to give sufficiently large  $|\gamma_{xxx}|$  and large electron correlation effects on  $\gamma_{xxx}$  values. The addition of tight *d* ( $\zeta_d=0.2$ ) function tends to slightly decrease the  $|\gamma_{xxx}|$  as well as the case of the neutral allyl (a). It is found that the electron correlation effects are mainly incorporated in the MP2 correction and higher-order correlation effects only provide small changes of  $\gamma_{xxx}$ .

The  $\gamma_{xxx}$  values of three charged states of allyl for each basis set are plotted in Figure 3. As shown in the graphs, the double zeta (DZ) type basis set augmented by diffuse *p* and *d* ( $\zeta_p=0.05$ ,  $\zeta_d=0.05$ ) and the MP2 correlation methods must be used at least to describe the relative tendencies of the systems. Particularly, both the diffuse *p* and *d* functions are essential for describing the qualitative tendency of the anion state.

The order in the magnitude of  $\gamma_{xxx}$  values of allyl systems is obtained as  $|\gamma(\text{anion})| \gg |\gamma(\text{neutral})| > |\gamma(\text{cation})|$ . It is shown that the signs of  $\gamma_{xxx}$  are positive for neutral (a) and anion (c) states of allyl systems, while the sign of  $\gamma_{xxx}$  is negative for allyl cation (b). In the next section, the  $\gamma$  values of alternant and condensed-ring conjugated systems are investigated by the use of the MP2 correlation methods using the DZ + diffuse *p* and *d* basis set.

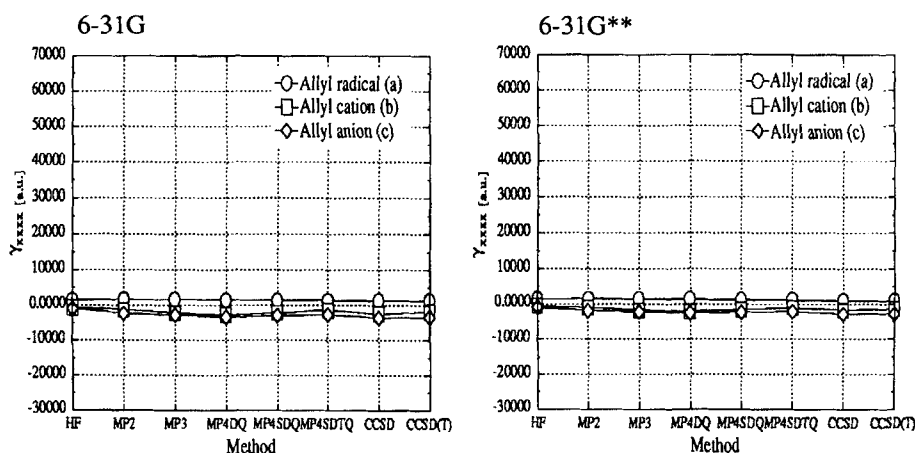


FIGURE 3 Electron correlation dependences of  $\gamma_{xxx}$  for three charged states (neutral (a), cation (b) and anion (c)) of allyl in each basis set.

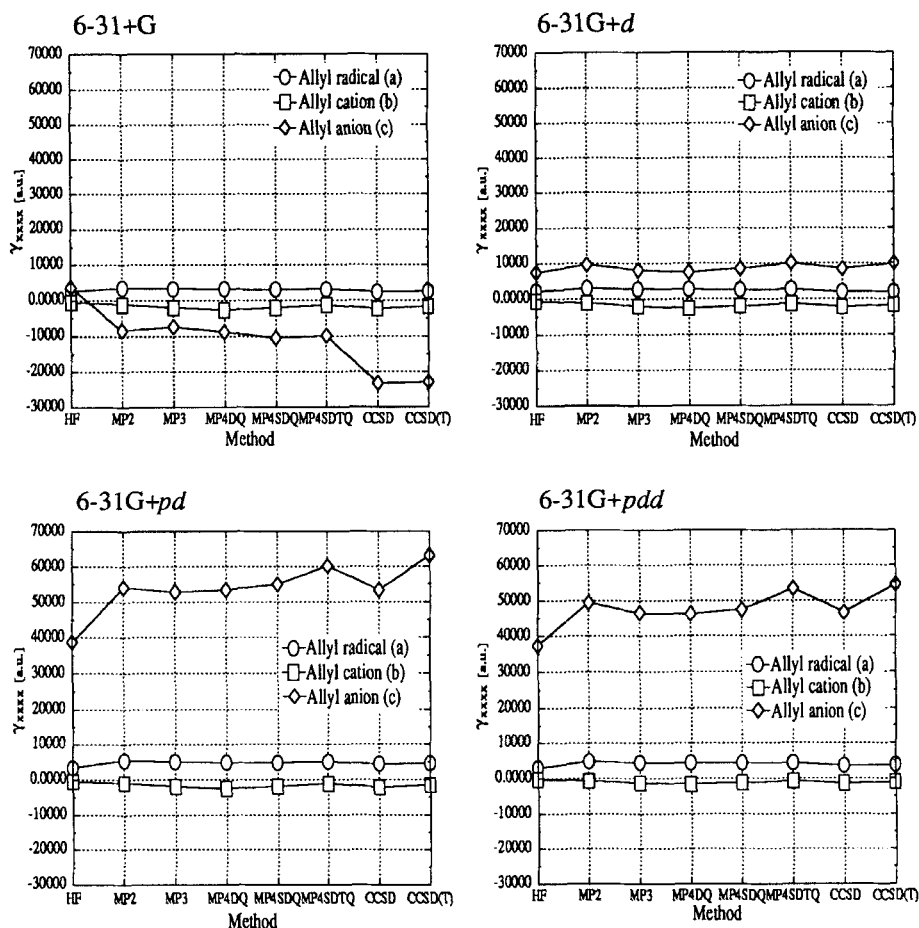


FIGURE 3 (Continued)

### MAIN COMPONENTS OF $\gamma$ FOR ALTERNANT AND CONDENSED-RING CONJUGATED SYSTEMS

Table I gives the static  $\gamma_{xxxx}$  (main component) values of triafulvalene (d) and pentalene (e) shown in Figure 4 calculated by CHF+MP2 method using 6-31G+*pd* ( $\zeta_p=0.05$ ,  $\zeta_d=0.05$ ) basis set. The geometries for three states (neutral, ion radical and ion states) of the systems are optimized by the HF method using 6-31G\*\* basis set.

The sign of  $\gamma_{xxxx}$  value of cation radical states of triafulvalene (d-2) as a model of alternant conjugated system is shown to be inverted between HF and MP2 methods. For the other charged states (d-1) and (d-3), it is found that the MP2 and HF methods



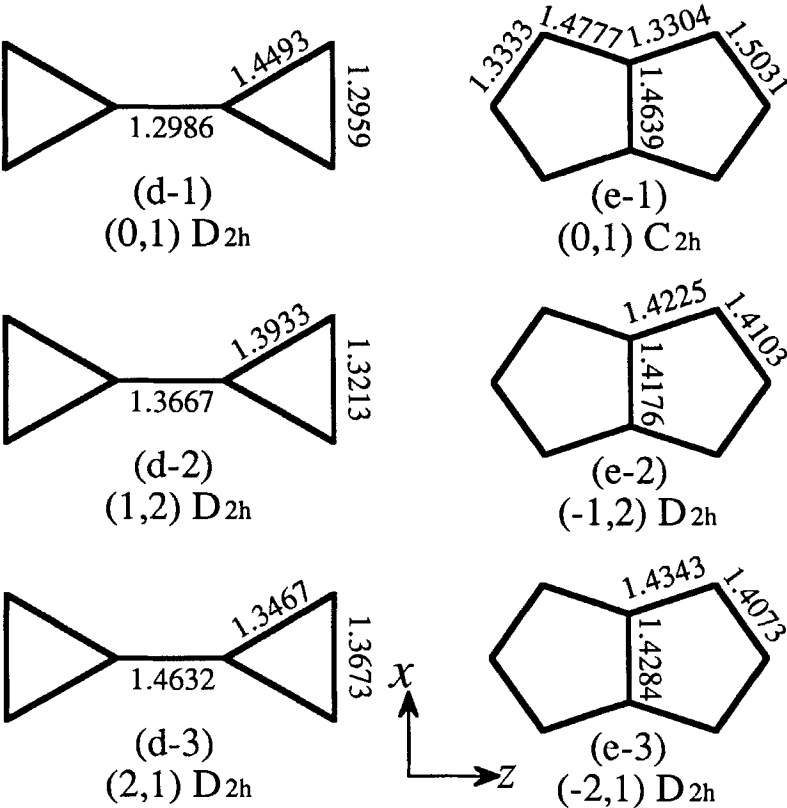


FIGURE 4 Calculated bond lengths [ $\text{\AA}$ ] of three charged states (neutral, ion radical and ion) of alternant and condensed-ring conjugated systems ((d) triafulvalene and (e) pentalene) using 6-31G\*\*. Numbers in parentheses depict charge and spin multiplicity, respectively.

TABLE I  $\gamma_{\text{int}}$  [a.u.] values of triafulvalene (d) and pentalene (e) calculated by the HF and MP2 methods using 6-31G+*pd* basis set.

System (charge, multiplicity)	Symmetry	HF	MP2
(d-1) (0,1)	$D_{2h}$	12000	17300
(d-2) (1,2)	$D_{2h}$	-12600	10600
(d-3) (2,1)	$D_{2h}$	866	1390
(e-1) (0,1)	$C_{2h}$	4350	9330
(e-2) (-1,2)	$D_{2h}$	-11700	-105000
(e-3) (-2,1)	$D_{2h}$	177000	329000

give positive  $\gamma_{xxx}$  values though the magnitudes of  $\gamma_{xxx}$  are much enhanced in the MP2 in contrast to the HF method. The order in the magnitudes of  $\gamma_{xxx}$  at the MP2 level is obtained as  $|\gamma(\text{neutral})| > |\gamma(\text{cation})| \gg |\gamma(\text{dication})|$  which shows that the systems with larger number of electrons possess larger  $|\gamma_{xxx}|$  values.

As for the anion radical states of pentalene (e-2), which is a model of condensed-ring conjugated systems, it is interesting that both the HF and MP2 methods provide negative  $\gamma_{xxx}$  values. Furthermore, the  $|\gamma_{xxx}|$  values are shown to be extremely enhanced at MP2 level against the HF level. The electron correlation effects are found to be also large for neutral (e-1) and dianion (e-3) states. The order in the magnitudes of  $\gamma_{xxx}$  is obtained as  $|\gamma(\text{dianion})| > |\gamma(\text{anion})| \gg |\gamma(\text{neutral})|$ . The relations between the magnitude of  $\gamma_{xxx}$  of pentalene (e) and the number of  $\pi$ -electrons are similar with those of triafulvalene (d).

## RESULTS AND DISCUSSION

On the basis of the three-type approximation, the centrosymmetric systems with larger ground state polarizability  $\alpha_0$  than the excited polarizability  $\alpha_n$  have a possibility of exhibiting larger negative  $\gamma$  values. In this study, the  $\gamma$  values of neutral, ion radical and ion states of alternant and condensed-ring conjugated systems were examined by the *ab initio* CHF + electron correlation methods using extended basis sets. From the results of neutral (a), cation (b) and anion (c) states of allyl systems, it is found that the double zeta (DZ) + diffuse *p* and *d* ( $\zeta_p=0.05$ ,  $\zeta_d=0.05$ ) basis sets are at least necessary for describing the relative tendencies of the three states of allyl systems (a)-(c). For the allyl cation (b), it is found that the electron correlation effects are very large and that the  $\gamma_{xxx}$  values come to be negative in sign. This seems to be caused by larger mobility of  $\pi$ -electrons of allyl cation (b) due to the resonating-valence states than that of neutral allyl (a), which tends to enhance the type (II) contribution ( $\gamma > 0$ ). On the other hand, for the anion state of allyl (c), the diffuse orbitals are used by the excessive electrons to enhance the polarization in the presence of the external electric field, so that the  $|\gamma|$  values are seen to be enlarged where the type (III) (including higher excited states) seems to become important.

From the results by the HF and MP2 calculations of alternant and condensed-ring conjugated systems (triafulvalene (d) and pentalene (e)), the electron correlation effects on  $\gamma_{xxx}$  (main component) values are found to be extremely large for ion radical systems (d-2) and (e-2); the sign of  $\gamma$  for cation radical state of triafulvalene (d-2) is inverted between HF and MP2 methods, while the magnitude of  $\gamma$  dramatically

increase for the anion radical states of pentalene (e-2). The behaviors of  $\gamma$  seem to be due to the mobility of electrons as well as the case of allyl cation (b). However, the  $\gamma_{\text{anion}}$  value of cation radical state of triafulvalene (d-2) is not negative in sign in contrast to that of anion radical state of pentalene (e-2). Therefore, we must perform the *ab initio* calculations with higher-order electron correlation corrections (beyond the second-order ones) in order to confirm the tendencies of  $\gamma$  for alternant and condensed-ring conjugated systems. Particularly, the condensed-ring systems are very interesting since they are involved in fullerenes. From the results obtained here, the poly anion radical states of fullerenes, e.g.,  $C_{60}^{\cdot-}$  and  $C_{70}^{\cdot-}$ , are expected to possess large negative  $\gamma$  values and exhibit large correlation effects on  $\gamma$  values. In order to confirm the predictions, *ab initio* calculations with the moderate-size extended basis sets (DZ + diffuse *p* and *d*) including the MP and CC correlations for large-size alternant and condensed-ring conjugated systems will be also performed in the next paper II.

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